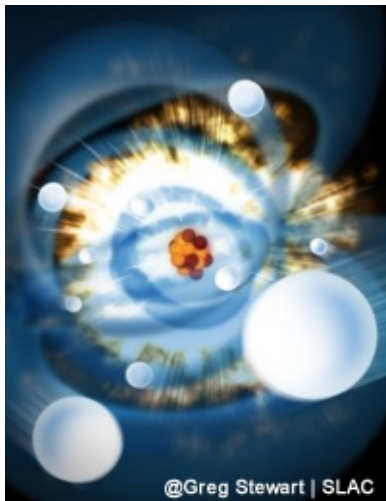


Non-resonant and resonant multiphoton processes with ultraintense x-rays

Linda Young



XFEL User Meeting 2011
Hamburg, Germany
26-28 January 2011



Outline

- Context & motivation
- Non-resonant high intensity x-ray phenomena
LCLS Experiment 1: Oct 1 - 6, 2009
- Resonant high intensity x-ray processes
LCLS Experiment 5: Oct 29 - Nov 3, 2009
- Summary



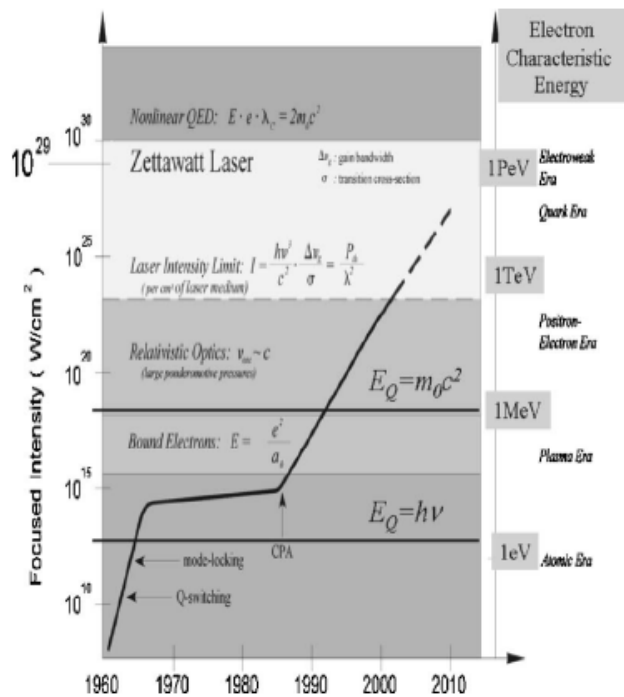
Compare the evolution of high intensity optical and x-ray sources

High-intensity at optical wavelengths

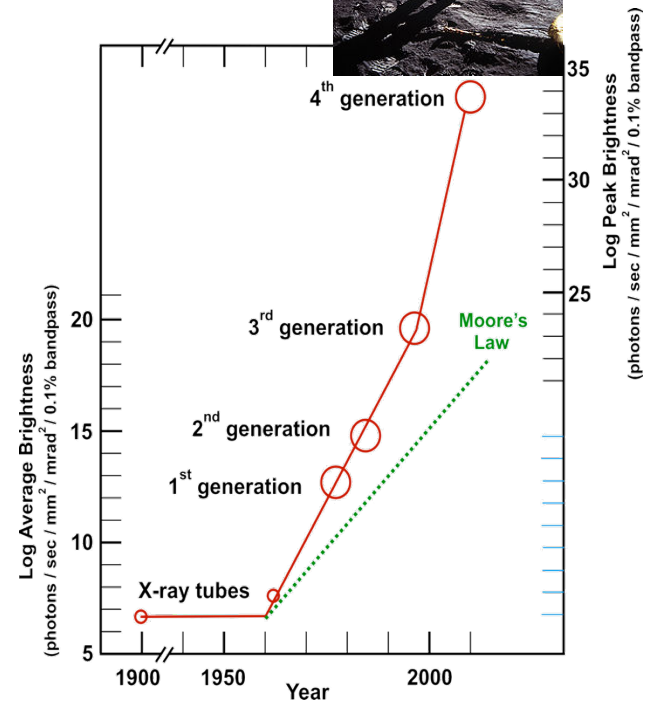
- high harmonic generation
- tabletop coherent x-ray radiation
- attosecond pulses

High-intensity at x-ray wavelengths

?
?
?



G. Mourou RMP 2006



D. Moncton, George Brown



Contrast optical and x-ray interactions at high intensity

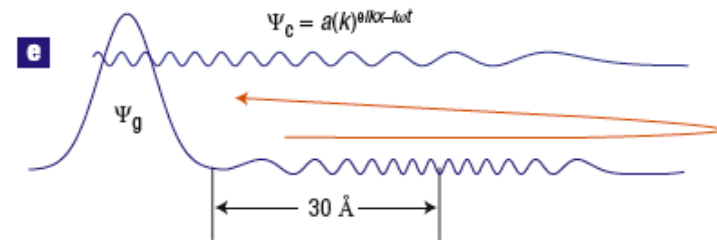
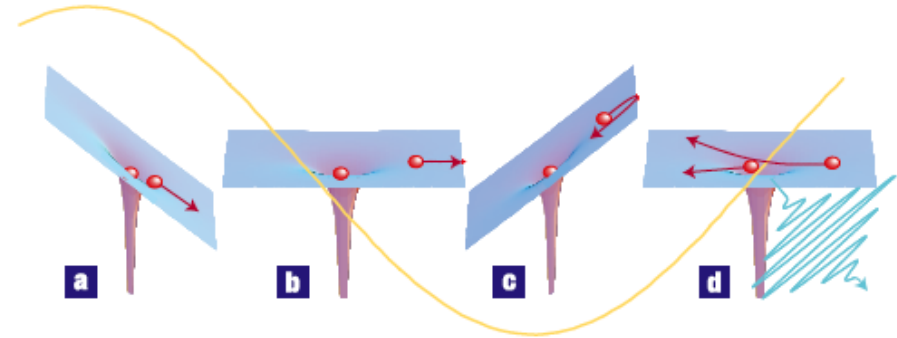
At long wavelengths - laser-driven electron dynamics is dominant
 ... not so at short wavelengths

electron ponderomotive energy (au)

$$U_p = I/4\omega^2$$

displacement

$$\alpha = E/\omega^2$$



Graphic from Corkum & Krausz
 Nature Physics (2007)

Ti:sapphire laser (1.55 eV) PW/cm²
 $U_p \sim 60 \text{ eV} \ \& \ \alpha \sim 50 \text{ au}$

LCLS (800 eV) 100 PW/cm²
 $U_p \sim 25 \text{ meV} \ \& \ \alpha \sim 0.003 \text{ au}$

Science Drivers for LCLS



AMO: Atomic Molecular and Optical

SXR: Soft X-ray Materials Science

XPP: X-ray Pump-Probe

XCS: X-ray Correlation Spectroscopy

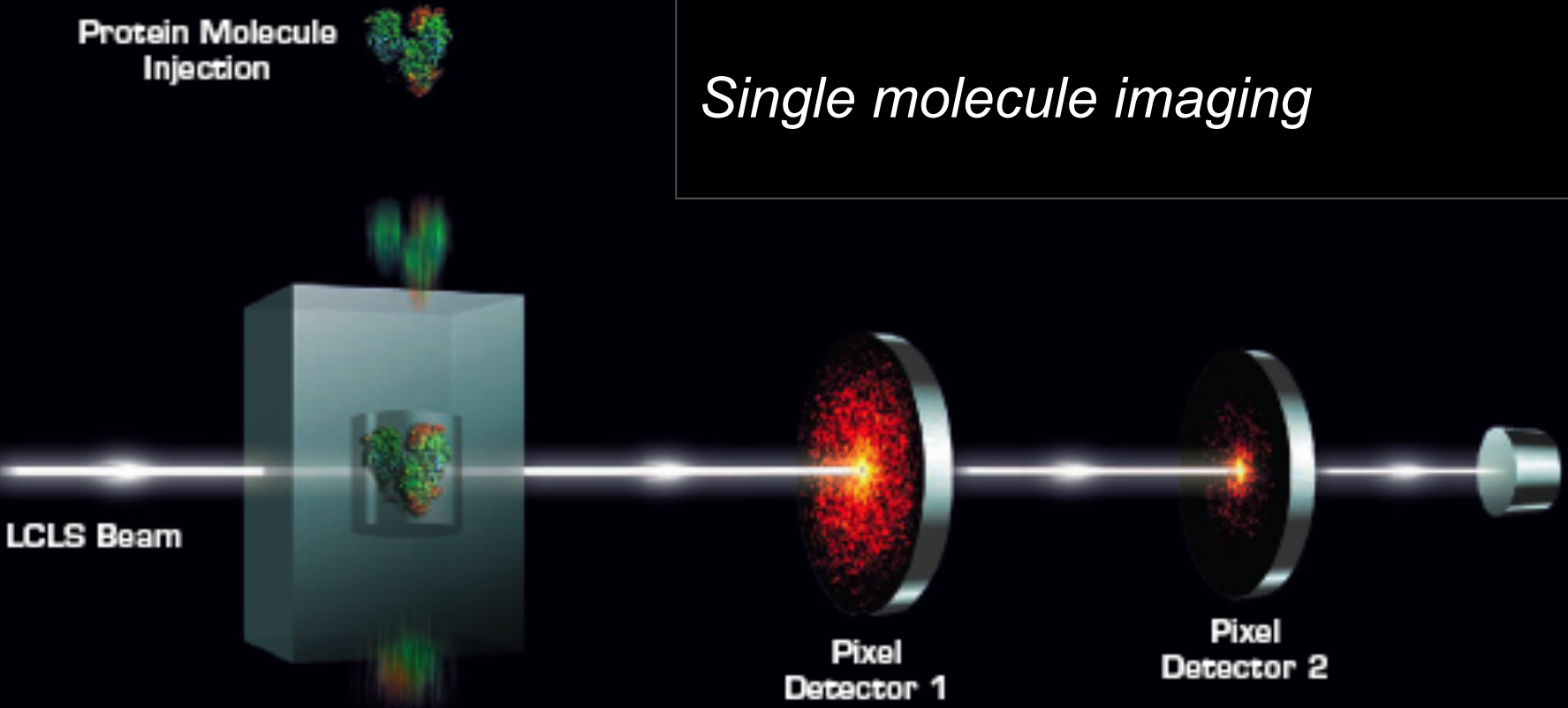
CXI: Coherent X-ray Imaging

MEC: Materials in Extreme Conditions

AMO

- Understand and control x-ray atom/molecule interactions at ultrahigh x-ray intensity as a foundation for other applications.
- Provide diagnostics of the LCLS radiation

Single molecule imaging



Protein Molecule Injection

LCLS Beam

Pixel Detector 1

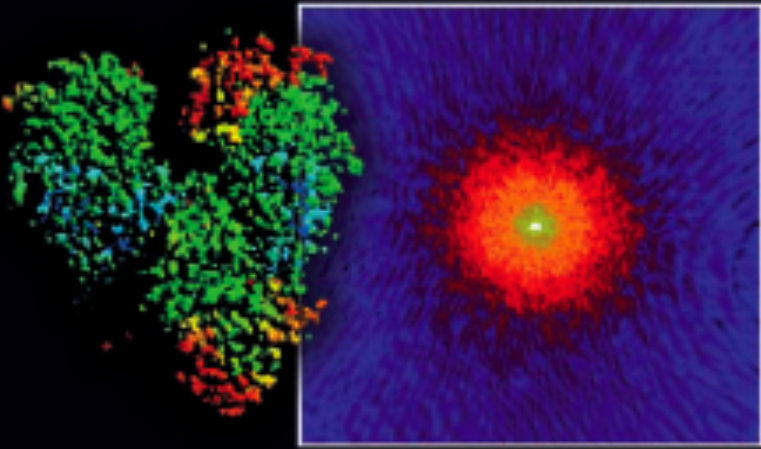
Pixel Detector 2

To Mass Spectrometer

Protein Molecule

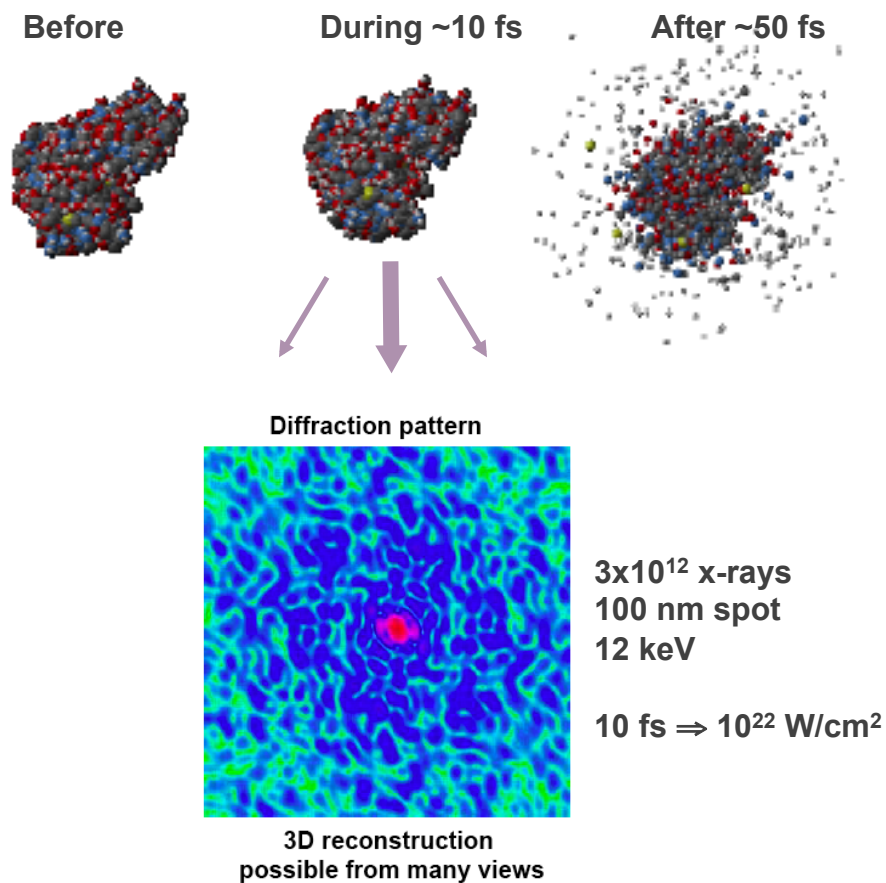
X-ray Diffraction Pattern

Operating with ultrafast pulses, LCLS will take images of molecules dropped into the x-ray beam. Scientists will merge the series of diffraction patterns of the molecules in many different positions. The resulting three-dimensional reconstruction will reveal the structures of proteins that cannot be crystallized and thus studied any other way.



AMO questions at the ultraintense x-ray frontier

- fundamental nature of x-ray damage at high intensity
 - Coulomb explosion
 - electronic damage
 - behavior at 10^{22} W/cm² - 1Å
- nonlinear x-ray processes
 - role of coherence
- quantum control of inner-shell processes



Neutze, Wouts, van der Spoel, Weckert, Hajdu Nature 406, 752 (2000)



LCLS Experiment 1 - Oct 1, 2009

Nature of the electronic response to

10^5 x-rays/ \AA^2
80 - 340 fs
800 - 2000 eV

$\sim 10^{18}$ W/cm²

Original single molecule imaging parameters, Neutze et al. Nature (2000)

3×10^{12} x-rays/(100 nm)² = 3×10^6 x-rays/ \AA^2

10 fs

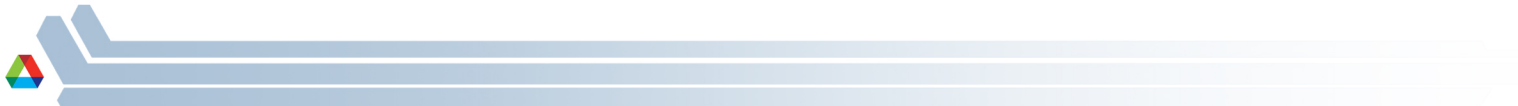
$\sim 10^{22}$ W/cm²



ARTICLES

Femtosecond electronic response of atoms to ultra-intense X-rays

L. Young¹, E. P. Kanter¹, B. Krässig¹, Y. Li¹, A. M. March¹, S. T. Pratt¹, R. Santra^{1,2}, S. H. Southworth¹, N. Rohringer³, L. F. DiMauro⁴, G. Doumy⁴, C. A. Roedig⁴, N. Berrah⁵, L. Fang⁵, M. Hoener^{5,6}, P. H. Bucksbaum⁷, J. P. Cryan⁷, S. Ghimire⁷, J. M. Glownia⁷, D. A. Reis⁷, J. D. Bozek⁸, C. Bostedt⁸ & M. Messerschmidt⁸



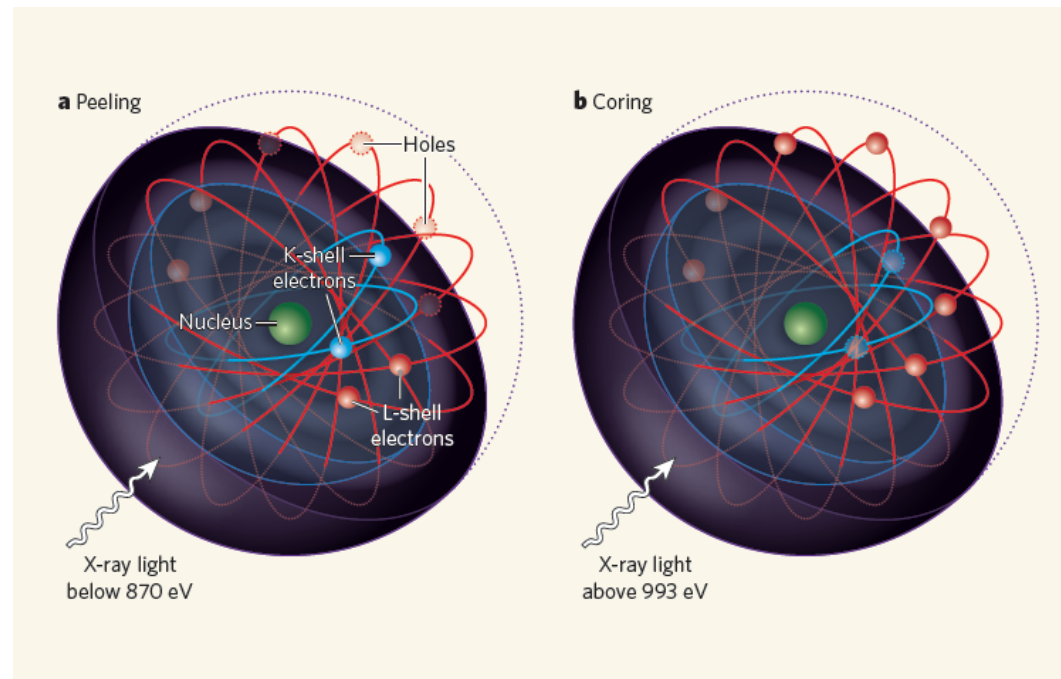
NEWS & VIEWS

ATOMIC PHYSICS

X-ray laser peels and cores atoms

Justin Wark

The world's first kiloelectronvolt X-ray laser produces such a high flux of photons that atoms can be 'cored'. In other words, the light source can knock out both the electrons of an atom's innermost shell.



Our approach to understanding ultraintense x-ray interactions

- Start with a well-characterized target

Binding energies in neutral neon

2p : ~21 eV

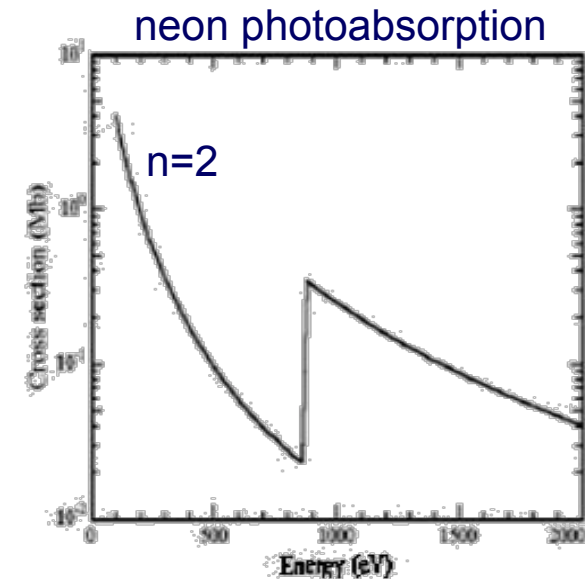
2s : ~48 eV

1s : ~870 eV

Inner-shell excitation

Auger yield 98%

Auger clock - τ_{1s} : 2.4 fs

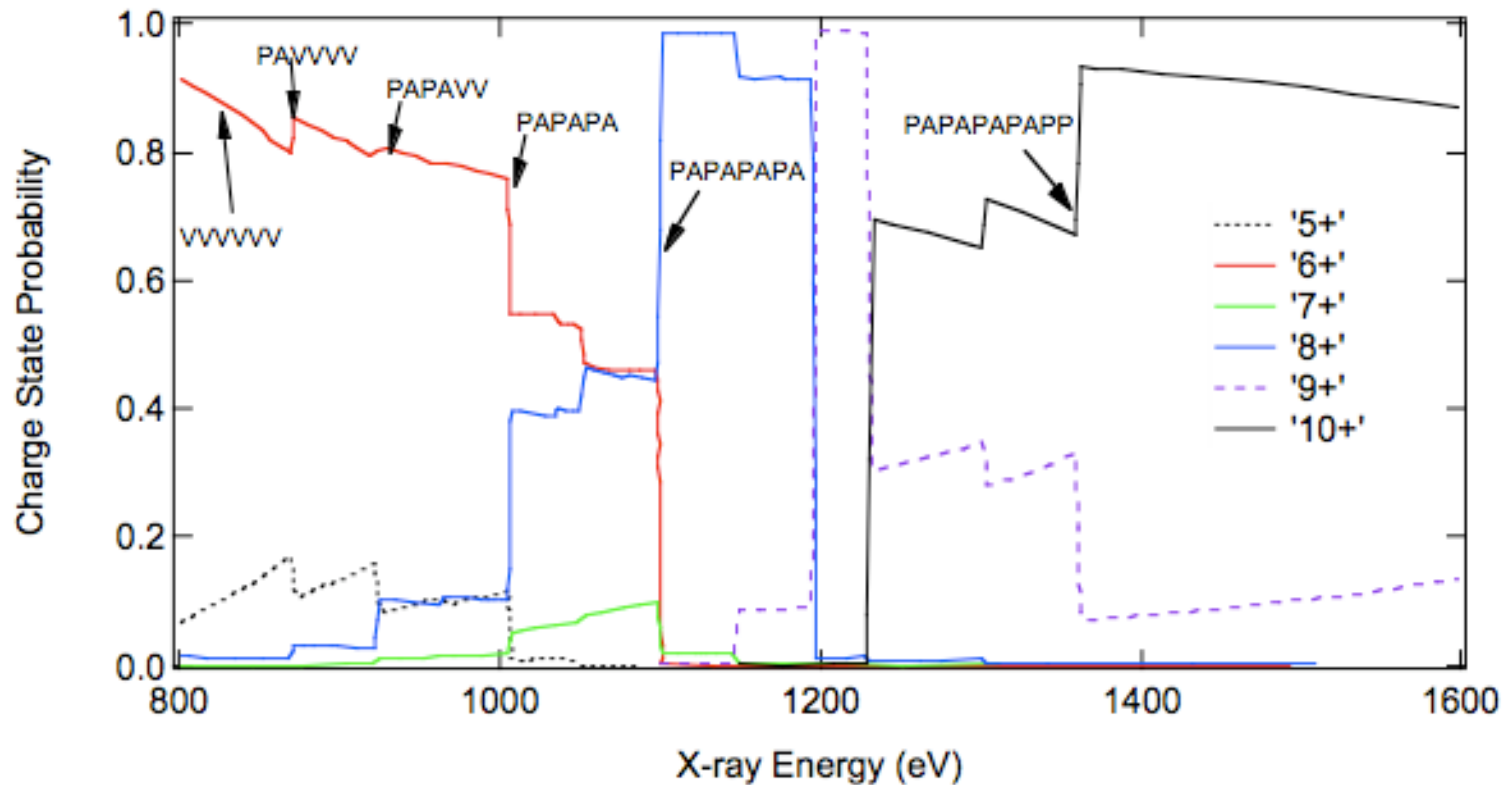


- Probe changes in interaction from outer- to inner-shell between 800-2000 eV



Guided by theory

Theory: Rohringer & Santra, PRA 76, 033416 (2007)

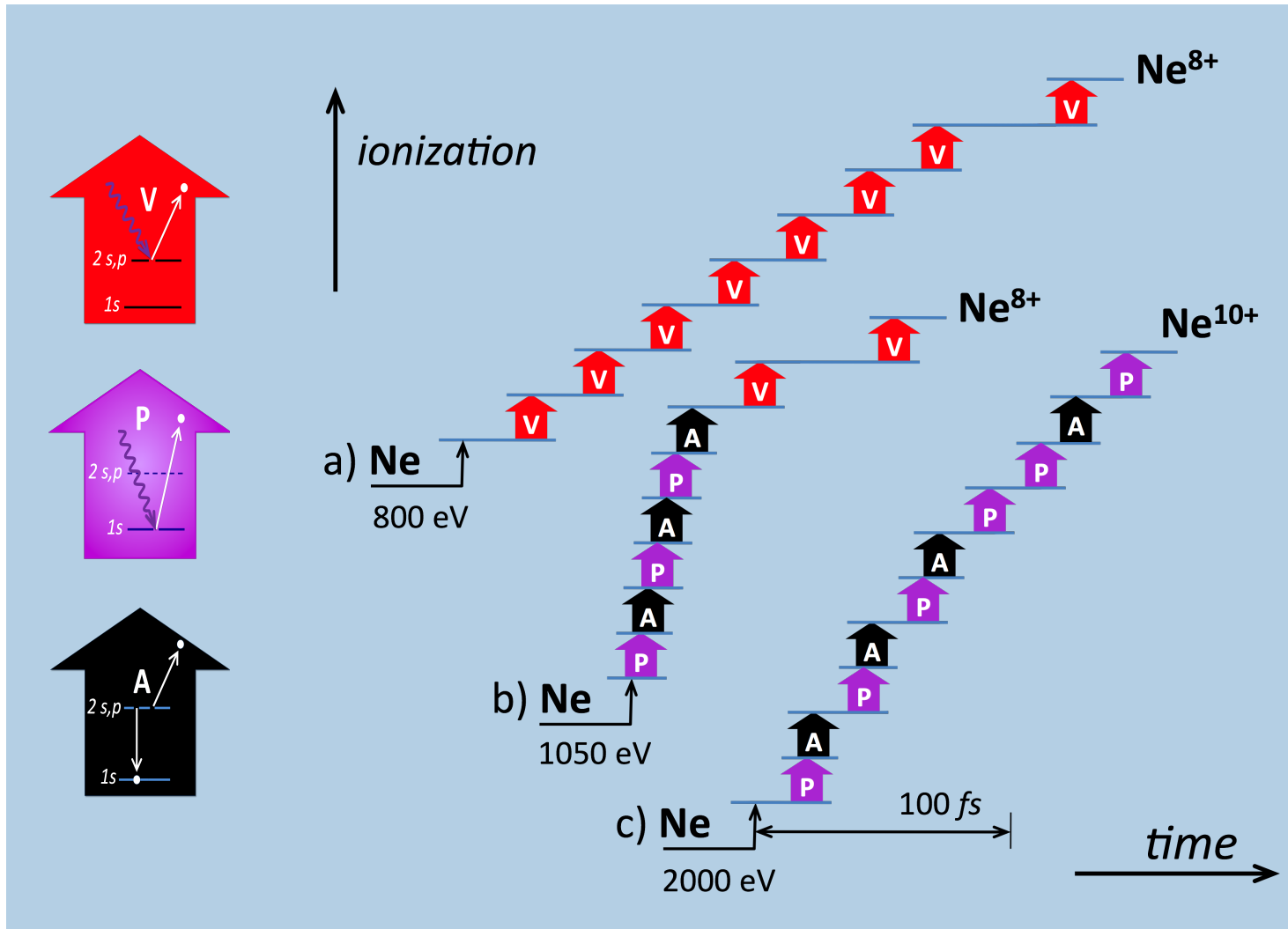


LCLS specs
10¹³ x-rays
230 fs
1 μm spot

Three target energies: 800 eV, 1050 eV, 2000 eV



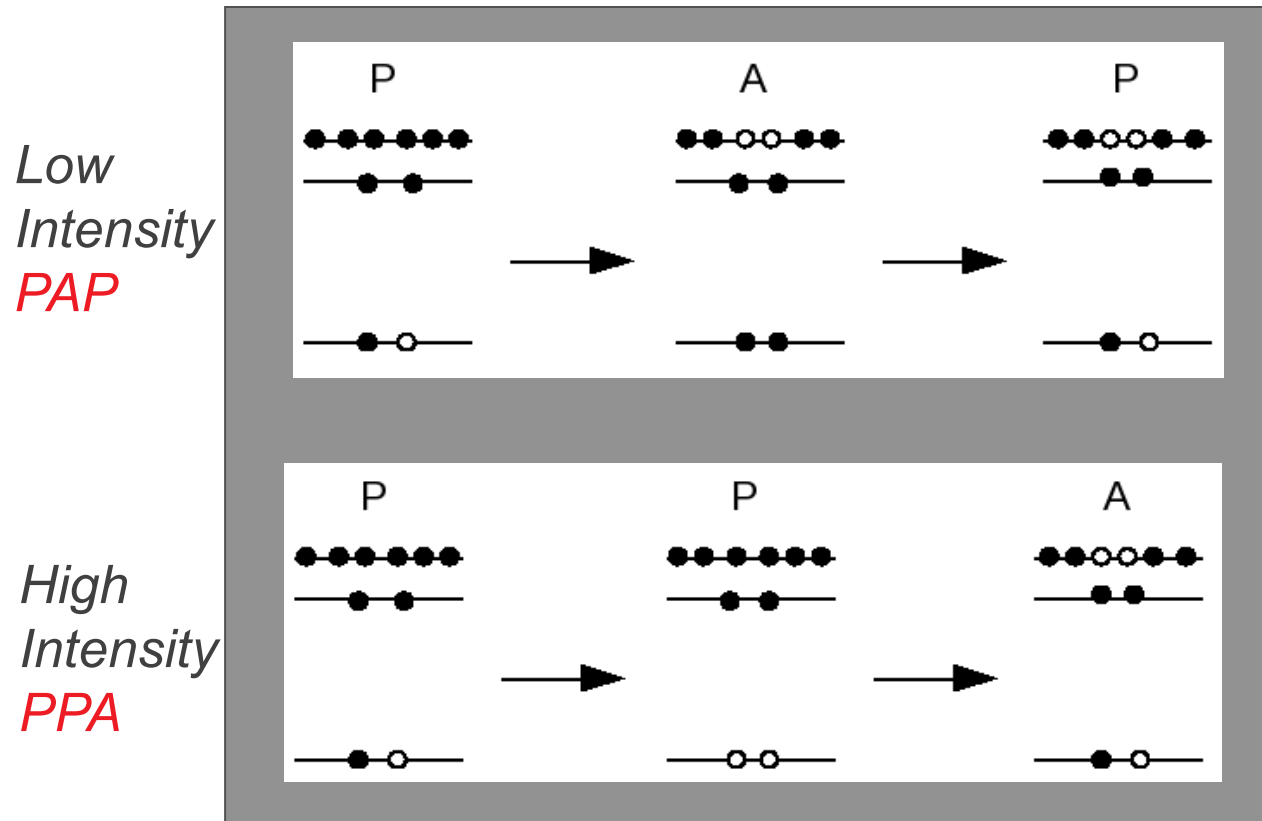
Valence ionization, core ionization and Auger decay



Sequential single photon processes dominate the interaction



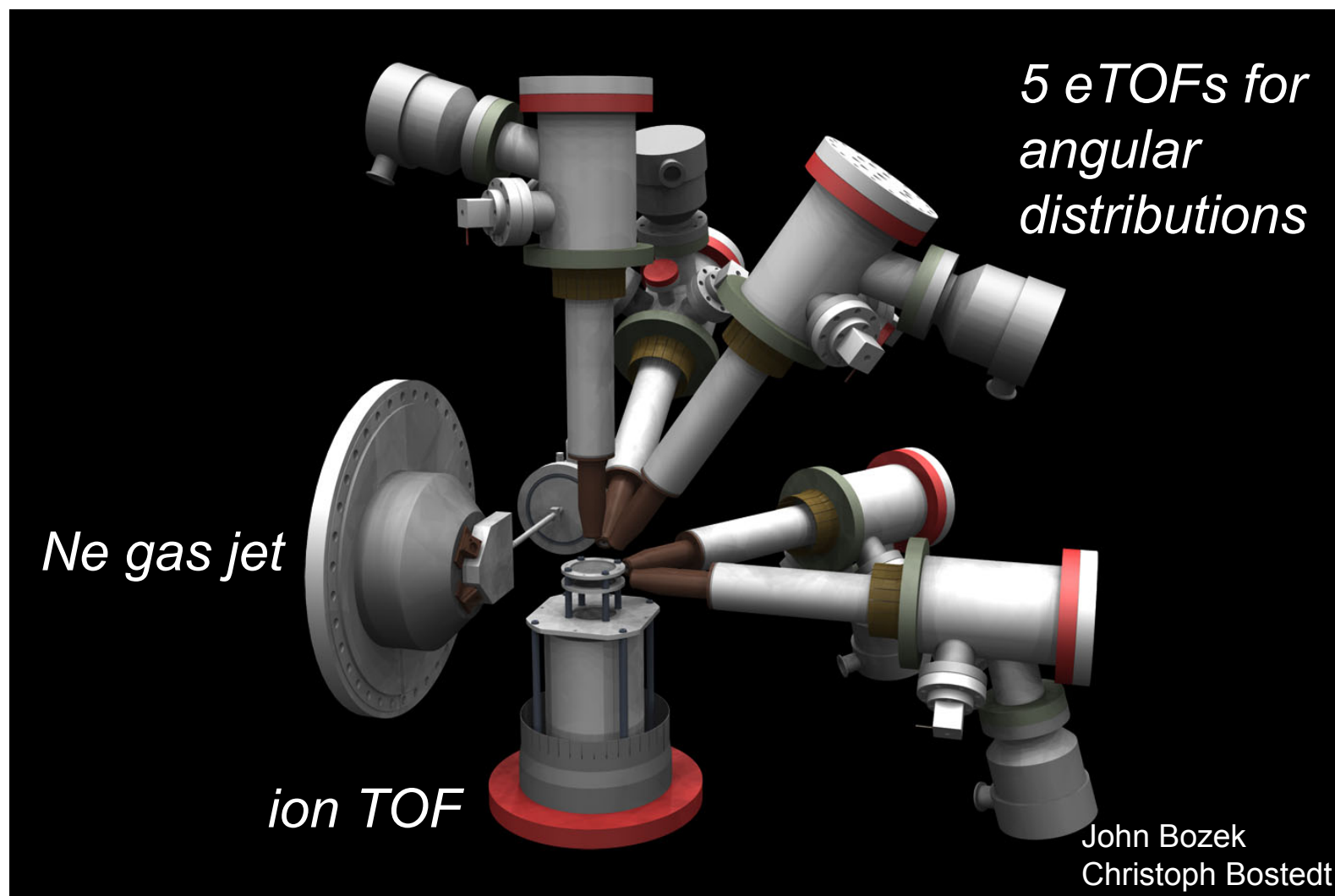
How does one arrive at a particular charge state?



- Hollow atoms produced at high x-ray intensity
- Electron spectroscopy can define the mechanism

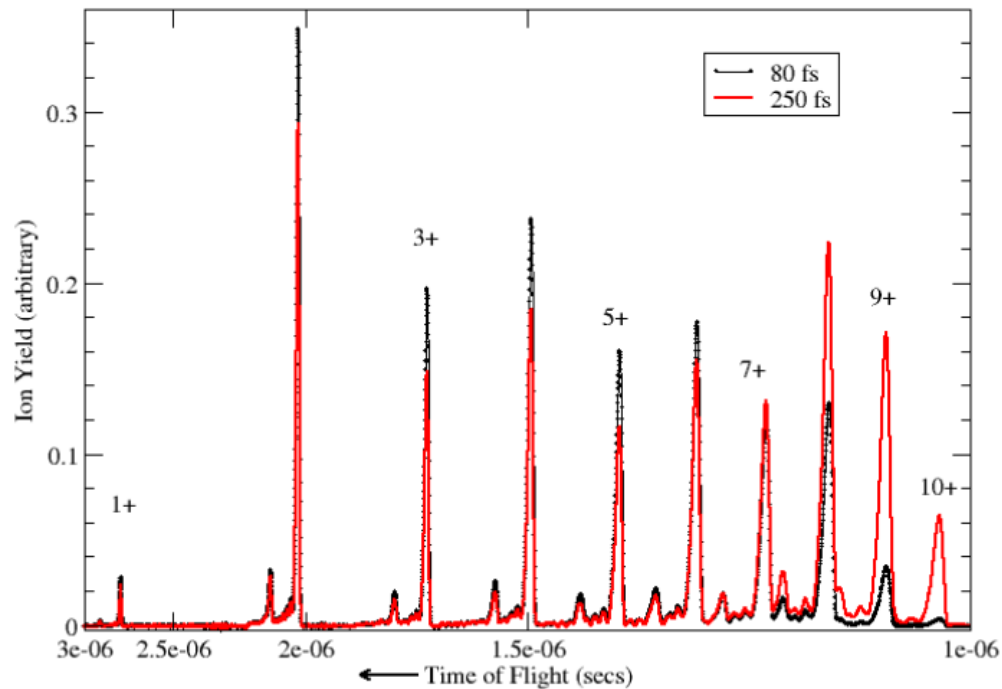


High field physics chamber



Day 1 - two interesting observations

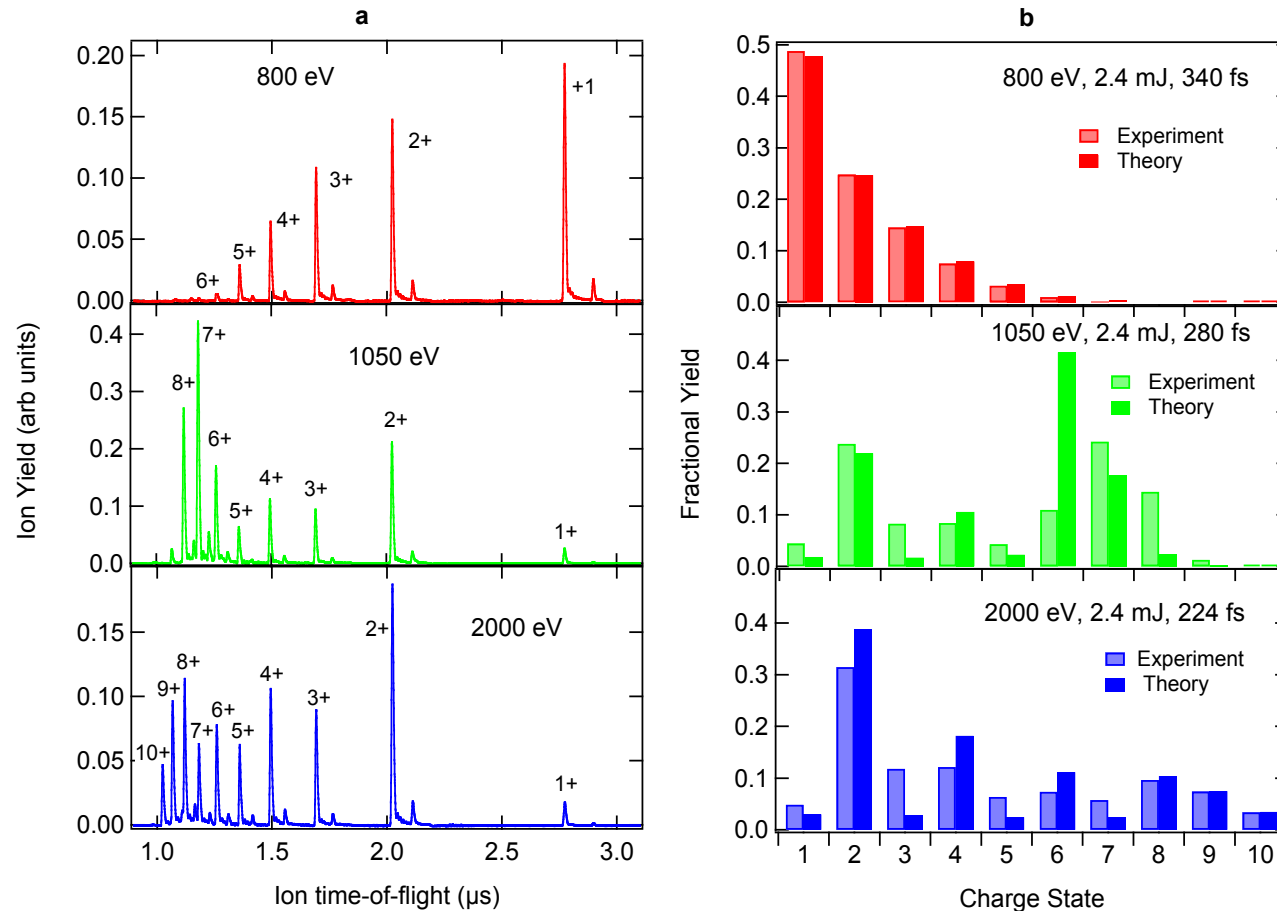
- Single ~ 100 fs pulse at 2000 eV fully strips neon
6-photon, 10-electron process



- Shorter pulses with equal pulse energy & fluence suppress absorption & damage.



Theory can model ultraintense x-ray-induced electronic damage in neon



Theory

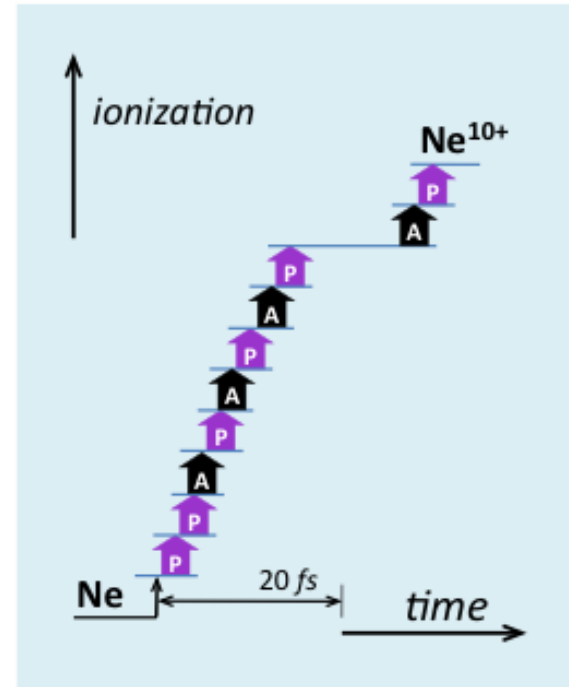
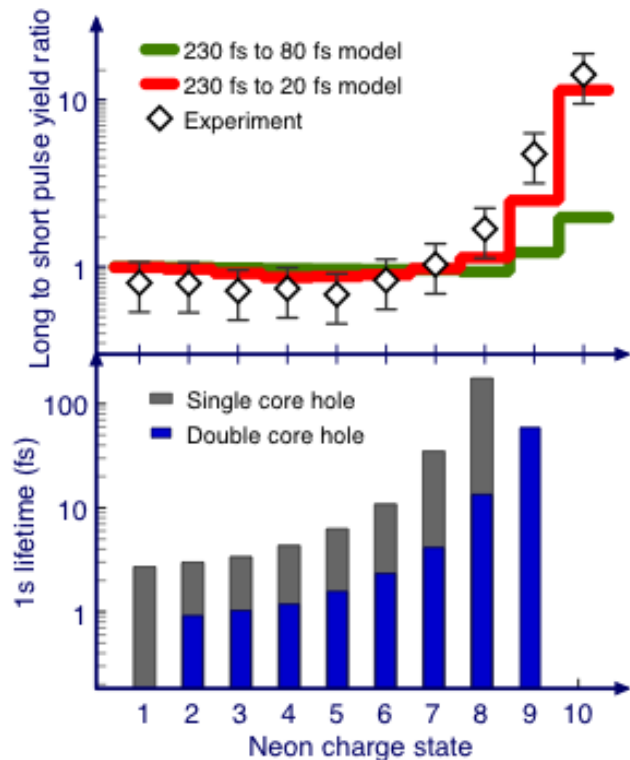
- Intensity averaged
- Fluence determined by experiment

Consistent with “measured” pulse energy and focus.

Sang-Kil Son, Robin Santra – refined calcs include shakeoff



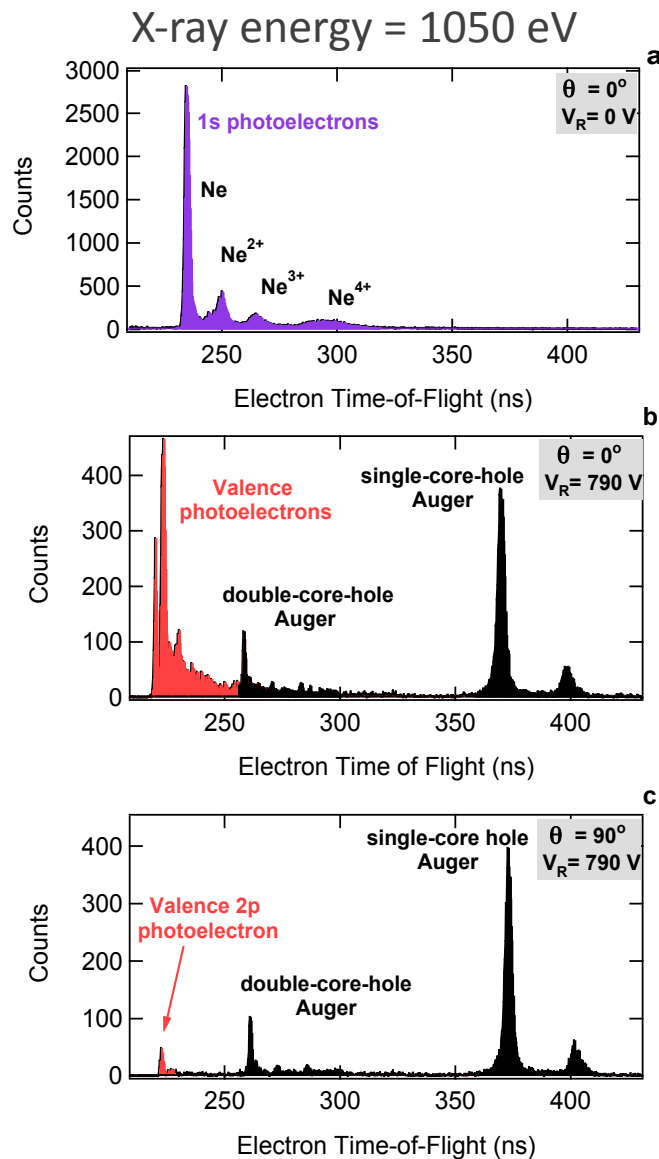
Atoms become transparent at high x-ray intensity !



- x-ray absorption is due to the presence of 1s electrons
- high x-ray intensities eject 1s electrons rendering the atom transiently transparent
- slowing atomic clocks create transparency at surprisingly long timescales



Electron spectrometers track ionization mechanism



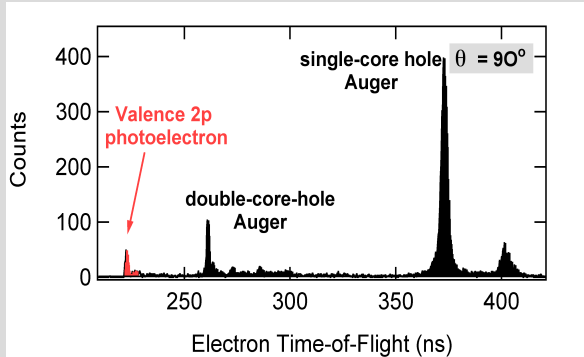
“Slow” 1s photoelectrons along x-ray polarization axis

“Fast” valence photoelectrons and Augers along polarization axis

Clean hollow atom signature
double-core-hole Auger
 $\theta = 90^\circ$

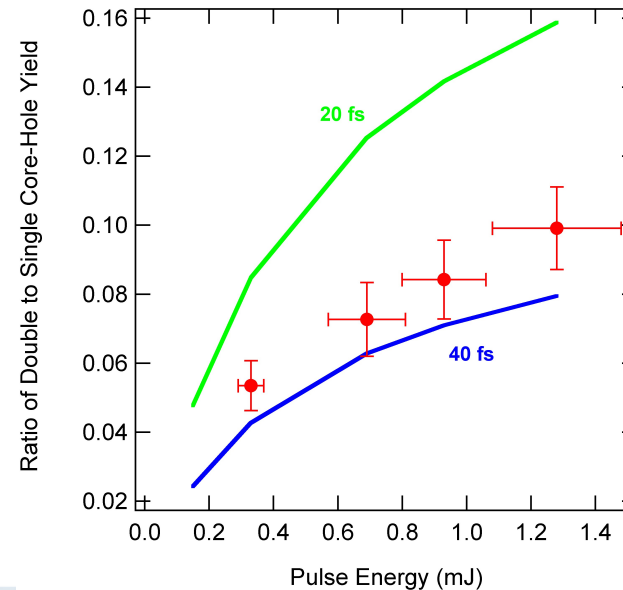


Hollow atom production: deliberate, huge and an an indicator of x-ray pulse duration

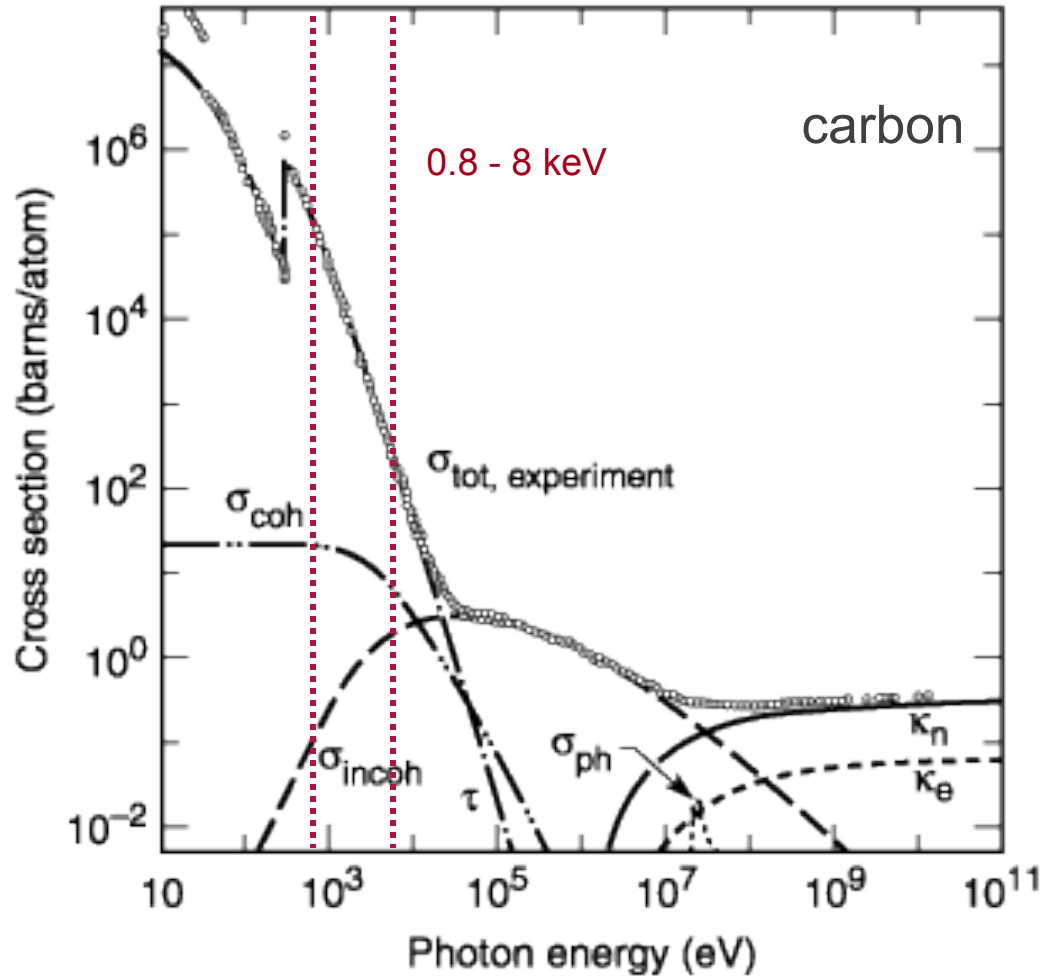


Hollow atom yield
@ LCLS ~10%
@ synchrotron ~0.3%
due to electron correlation

1050 eV,
nominal electron bunch
duration ~80 fs



Absorption vs scattering: normal and hollow atoms



	$\frac{\sigma_{\text{photo}}}{\sigma_{\text{elastic}}}$	$\frac{\sigma_{\text{Compton}}}{\sigma_{\text{elastic}}}$
2 keV	360	0.05
8 keV	20	0.60
8 keV hollow	2	

Impact of hollow atom formation on coherent x-ray scattering
 Sang-Kil Son, LY, RS
 Phys. Rev A. in press



Summary of ultra-intense x-ray interaction phenomena

- Target changes during a single 100 fs x-ray pulse at fluences similar to that for single molecule imaging
 - six-photon, ten-electron stripping of neon ($\sim 10^{12}/\mu\text{m}^2$)
 - multiphoton absorption probability high when fluence $> 1/\sigma$
- Intensity-induced x-ray transparency – a general phenomena
 - transient x-ray transparency caused by formation of hollow atoms
 - hollow atoms $\sigma_{\text{scatt}}/\sigma_{\text{abs}}$ is increased – advantageous for imaging
- Straightforward rate equation calculations capture essential physics
- Femtosecond time-scale atomic processes provide FEL diagnostics





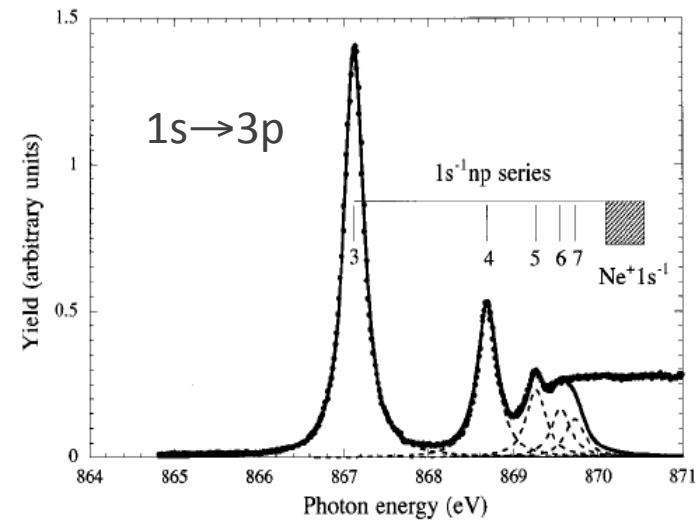
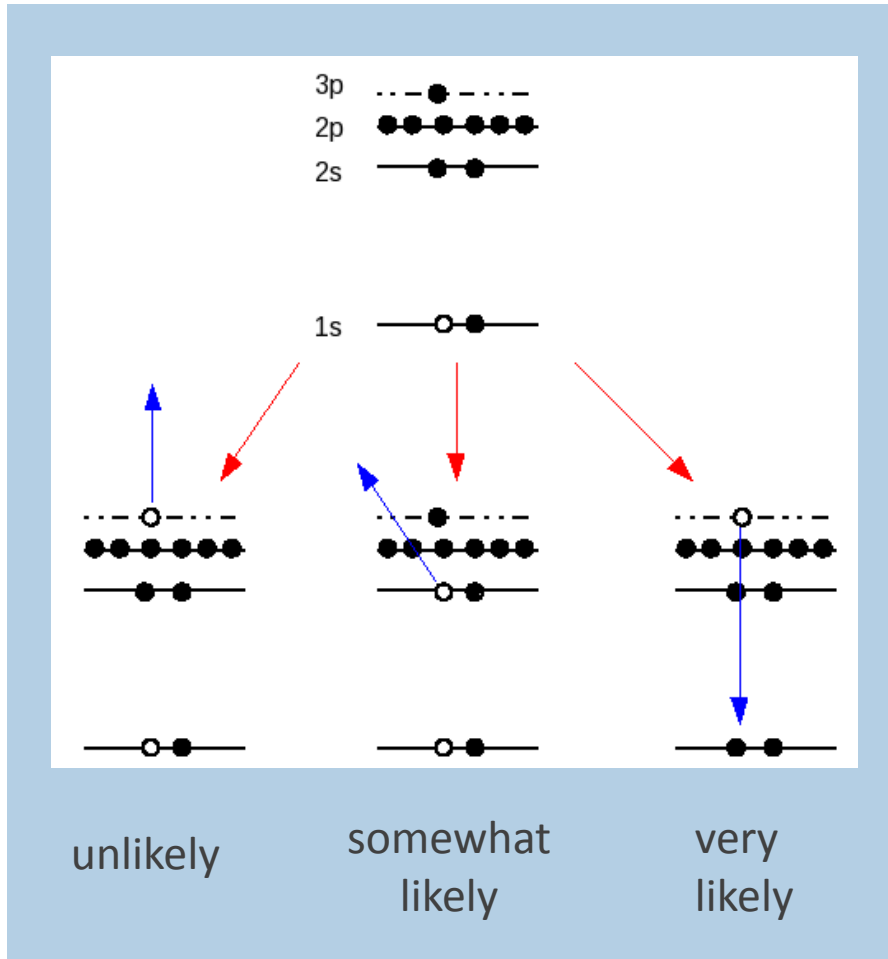
LCLS Experiment 5

Resonant x-ray processes at high intensity



Can we control inner-shell electron dynamics?

“Rabi flopping” may inhibit Auger decay & x-ray damage.



- Strong $1s \rightarrow 3p$ resonance

$$\mu_{Ne\ 1s-3p} = 0.01\ ea_0$$

$$\tau_{Ne\ 1s^{-1}} = 2.4\ fs = 100\ a.u.$$

- Rabi flopping possible

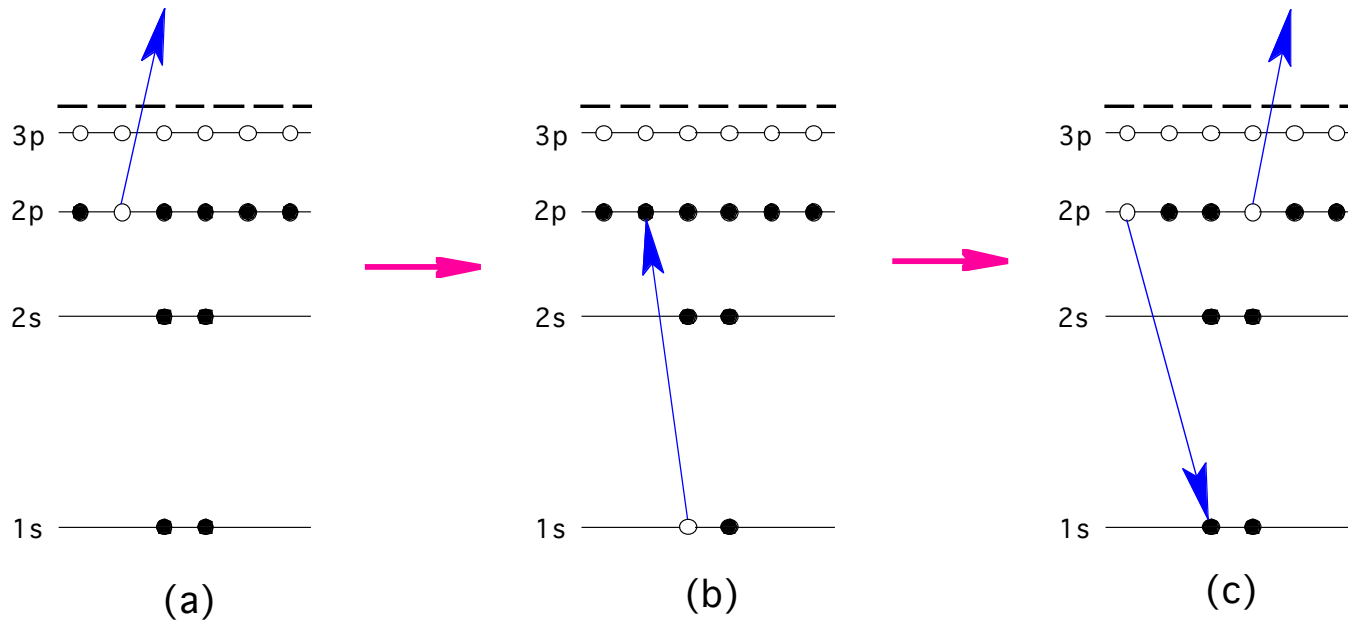
$$E_{Ne} \sim 6.3\ a.u.$$

$$I_{Ne} \sim 1.4 \times 10^{18}\ W/cm^2$$

But LCLS linewidth $\sim 8\ eV!$



Rabi-flopping on 1s - 2p resonance more feasible

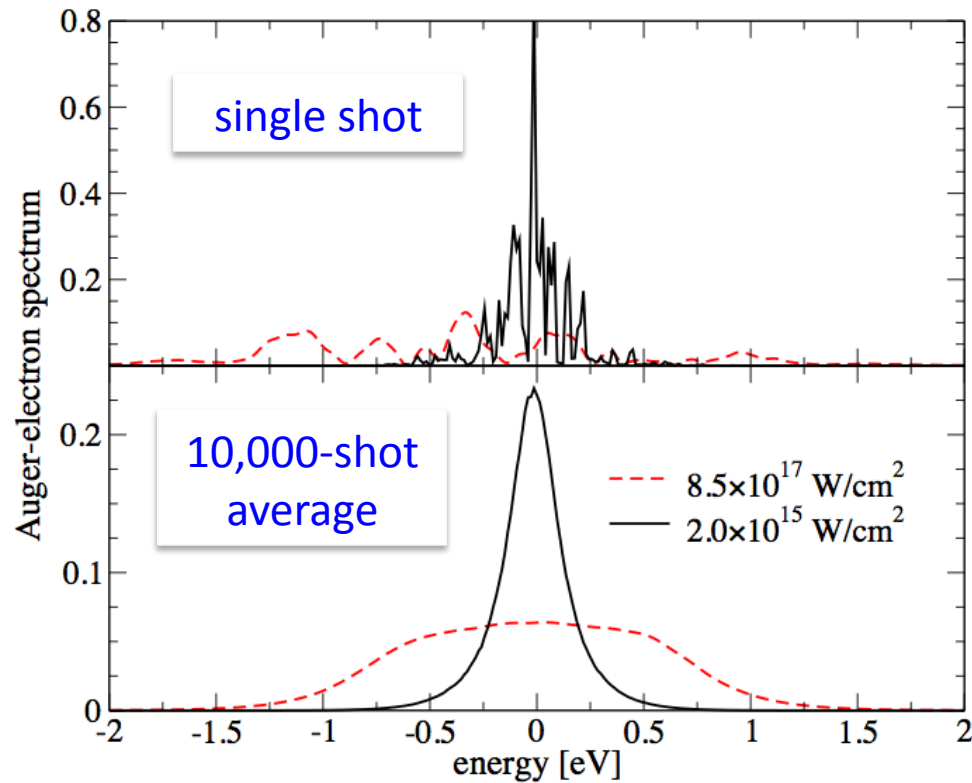


$$E_{x\text{-ray}} = 848.6 \text{ eV}$$

$$\sigma_{1s-2p} = 500\sigma_{2p-\infty} = 30 \sigma_{1s-3p}$$

Observe Auger yield when x-rays scanned over 1s - 2p resonance.
 Observe broadening at resonance to indicate Rabi flopping
 Theory: Rohringer & Santra PRA (2008).

Calculated “Resonant Auger effect at high x-ray intensity”



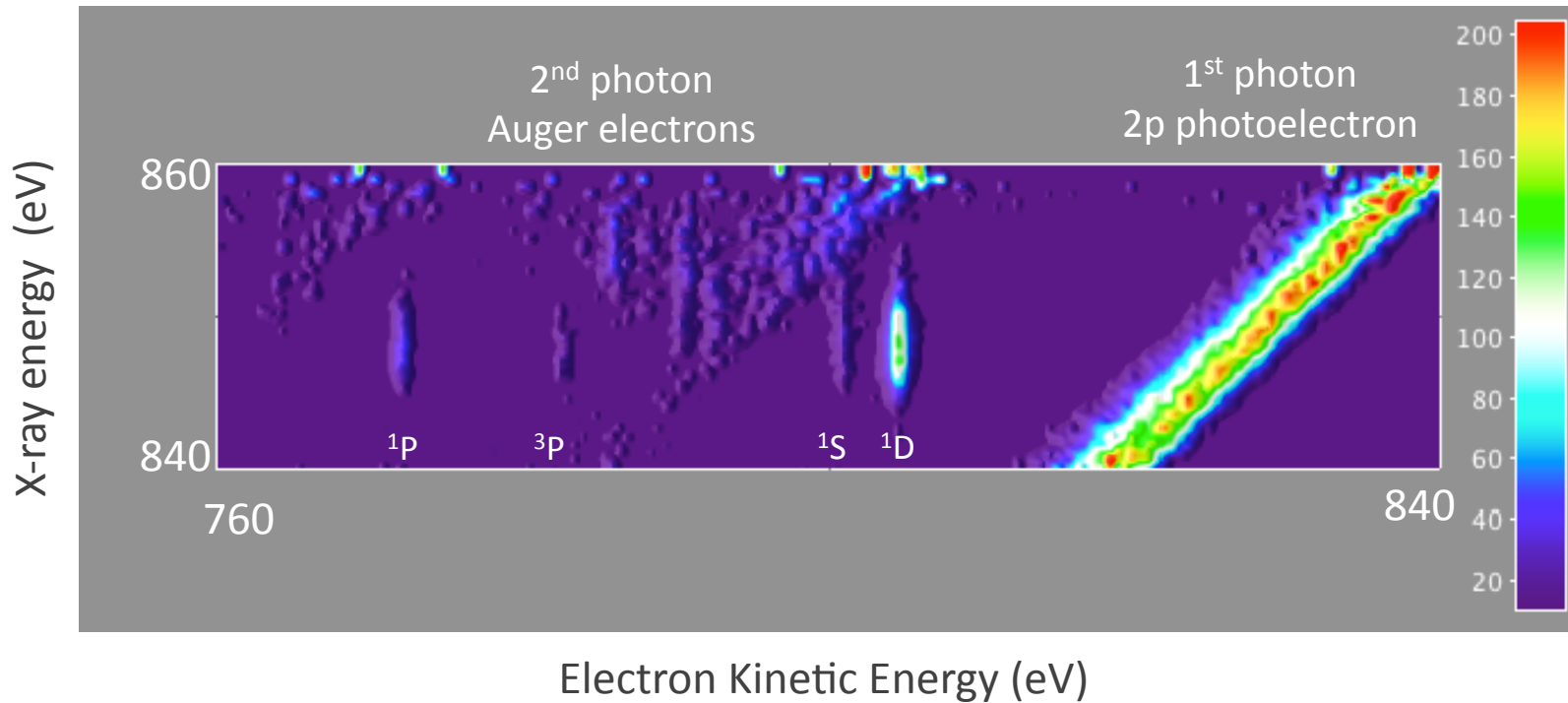
-> Look for Auger line broadening on resonance

N. Rohringer & R. Santra, PRA 77, 053404 (2008)



Electron spectra vs photon energy

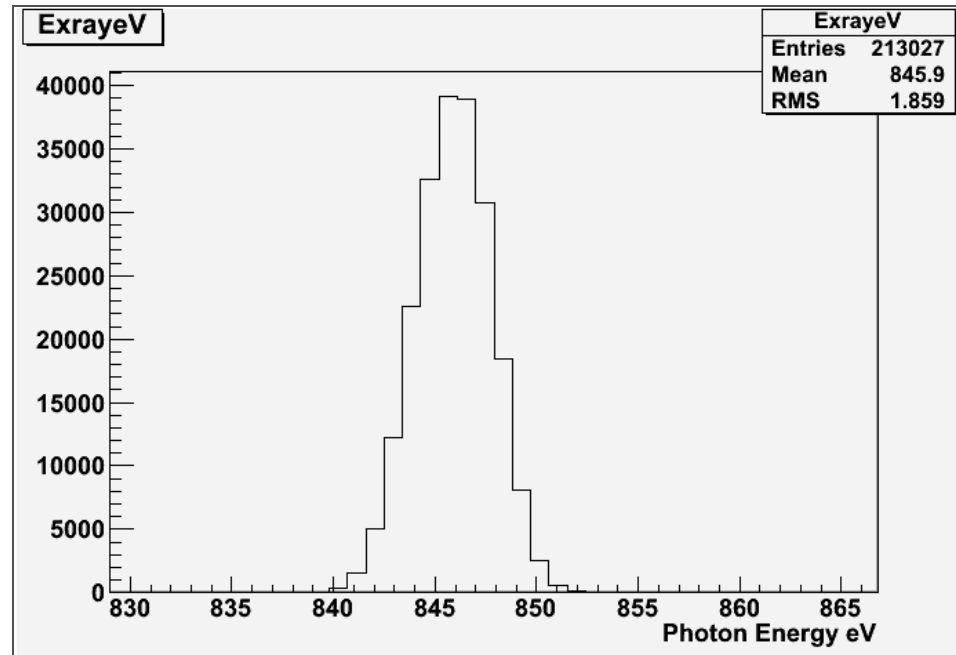
using eTOF1 perpendicular to photon polarization to suppress 2s photoelectrons (40 pC/bunch)



Preliminary from Bertold Krässig

Shot-to-shot photon energy jitter

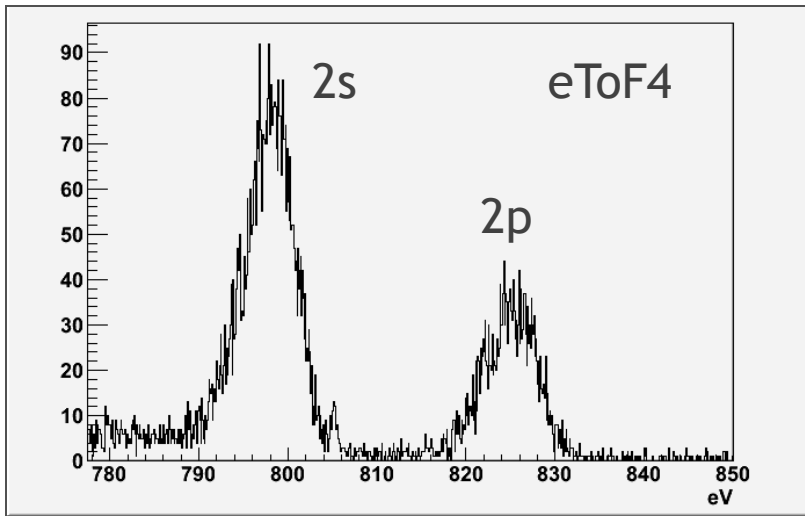
derived from GeV electron bunch energy measurement



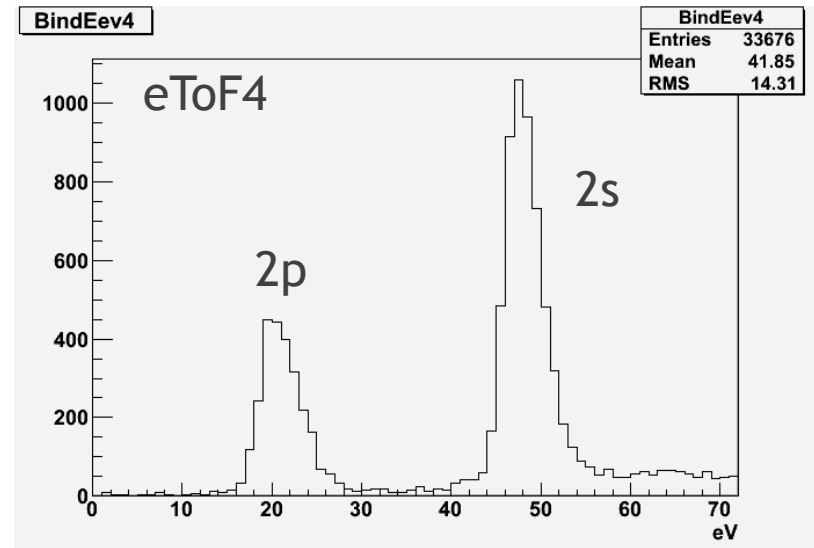
Conditions	FWHM photon energy jitter (eV)
40 pC (<10 fs) 850 eV 0.3 mJ 4500 A	4.25
250 pC (100 fs) 787 eV 1.5 mJ 2500 A	4.79
250 pC (100 fs) 769 eV 1.5 mJ 2500 A	5.24



Determination of intrinsic x-ray bandwidth from electron spectra



Electron kinetic energy (eV)

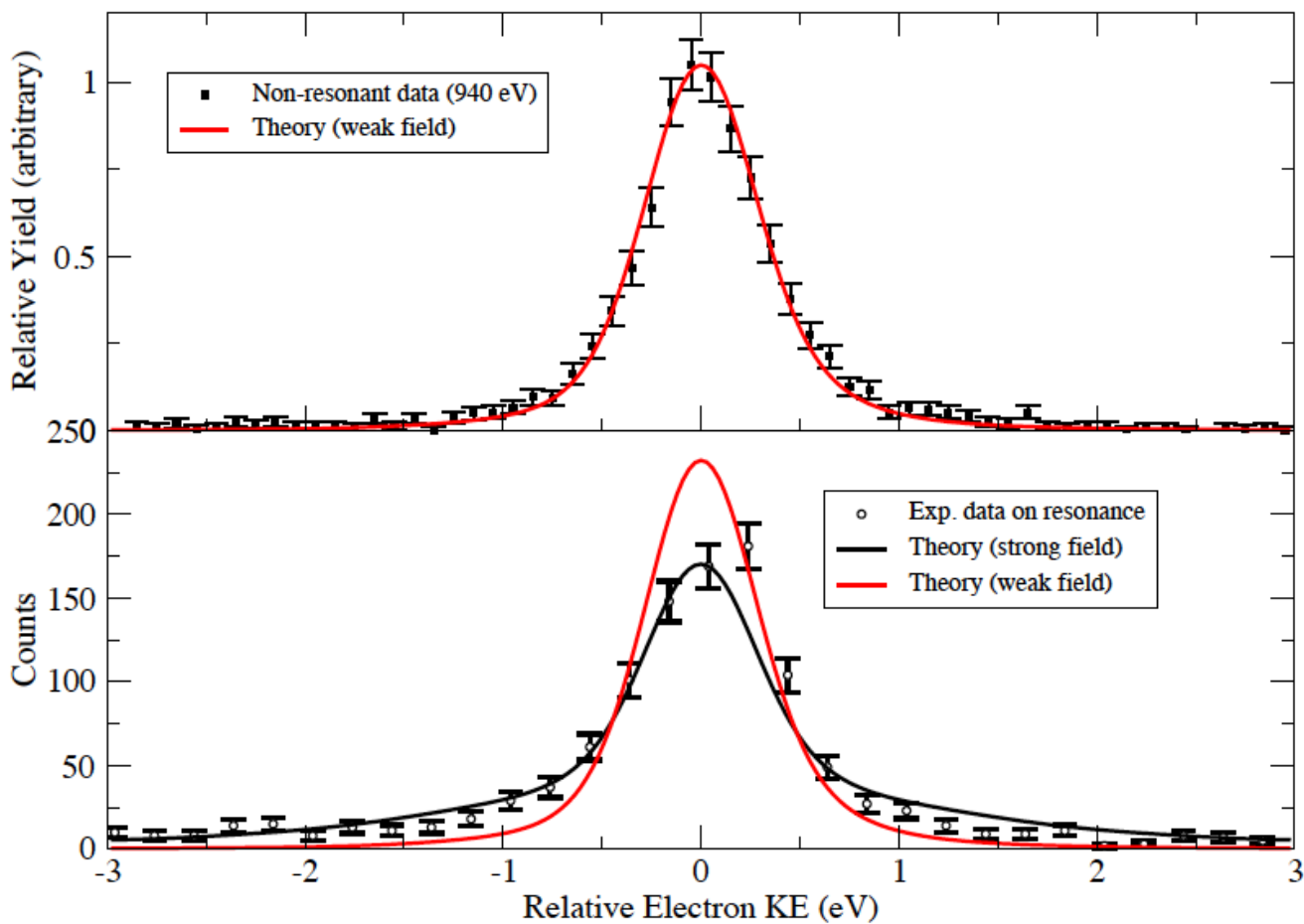


Electron binding energy (eV)

Conditions	Intrinsic x-ray pulse bandwidth (from 2s photopeak) (eV) (FWHM)	Intrinsic x-ray pulse bandwidth (from 2p photopeak) (eV) (FWHM)	Average bandwidth (eV) (FWHM)	%
40 pC (<10 fs) 850 eV 0.3 mJ 4500 A	4.3	4.5	4.4	0.5 %
250 pC (100 fs) 787 eV 1.5 mJ 2500 A	7.1	7.8	7.45	0.9%
250 pC (100 fs) 769 eV 1.5 mJ 2500 A	7.7	7.8	7.77	1%



Is the 1D Auger line broadened on 1s-2p resonance?



X-ray parameters
0.3 mJ, 8.5 fs, $2 \mu\text{m}^2$
(20% transmission)

Theory from N. Rohringer and R. Santra
Preliminary Analysis – E. Kanter



Summary

- Insight into ultraintense x-ray interactions
 - six-photon, ten-electron stripping of neon ($\sim 10^{12}/\mu\text{m}^2$)
 - multiple photon absorption probability high when fluence $> 1/\sigma$
- Intensity-induced x-ray transparency – a general phenomena
 - transient x-ray transparency caused by formation of hollow atoms
 - hollow atoms $\sigma_{\text{scatt}}/\sigma_{\text{abs}}$ is increased
- Femtosecond time-scale atomic processes provide FEL diagnostics
- Straightforward rate equation calculations capture essential physics
- Intense x-rays can “control” inner-shell electron dynamics

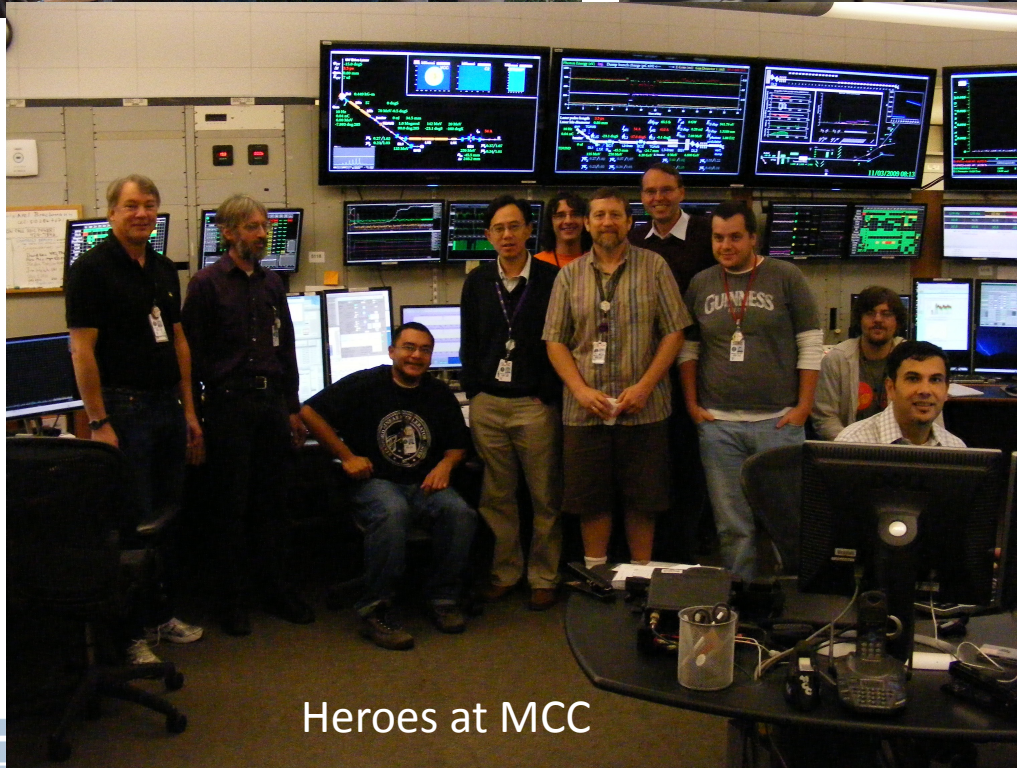




Argonne AMO group Oct 2009



Heroes at AMO Control



Heroes at MCC

